

Geislunar mynstur - skilar

Hertz-tvískauts geislun fjörsíð

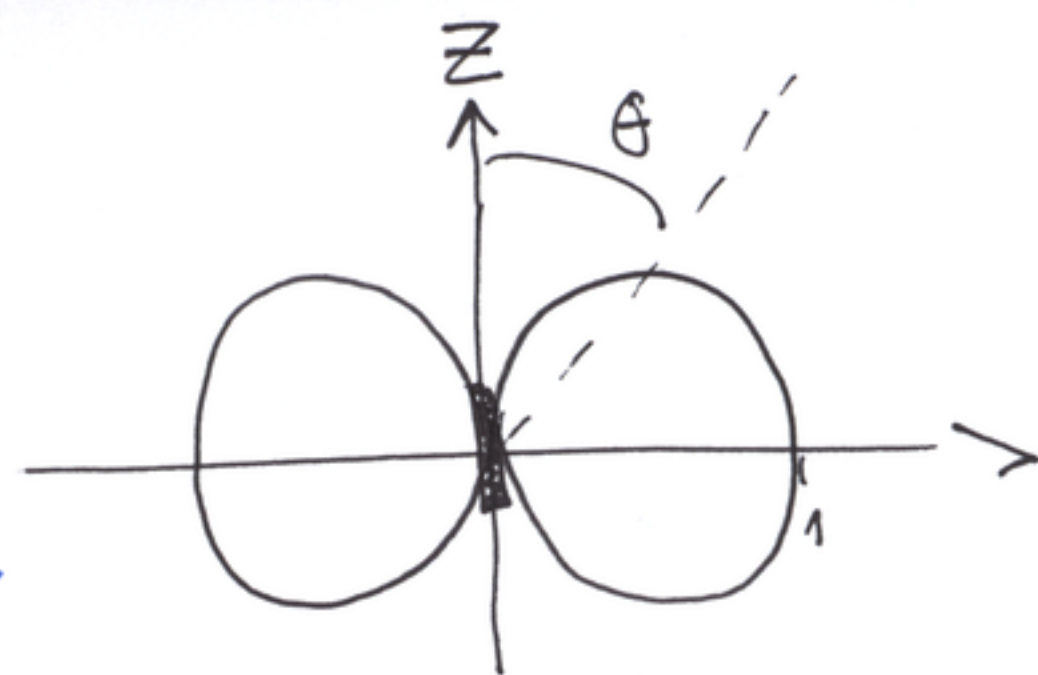
$$E_{\theta} \sim \frac{e^{-i\beta R}}{R} \sin \theta$$

$$H_{\phi} \sim \frac{e^{-i\beta R}}{R} \sin \theta$$

$$E_{\theta} \sim H_{\phi}, \text{ öðað } \phi$$

(E-plane): gefið R

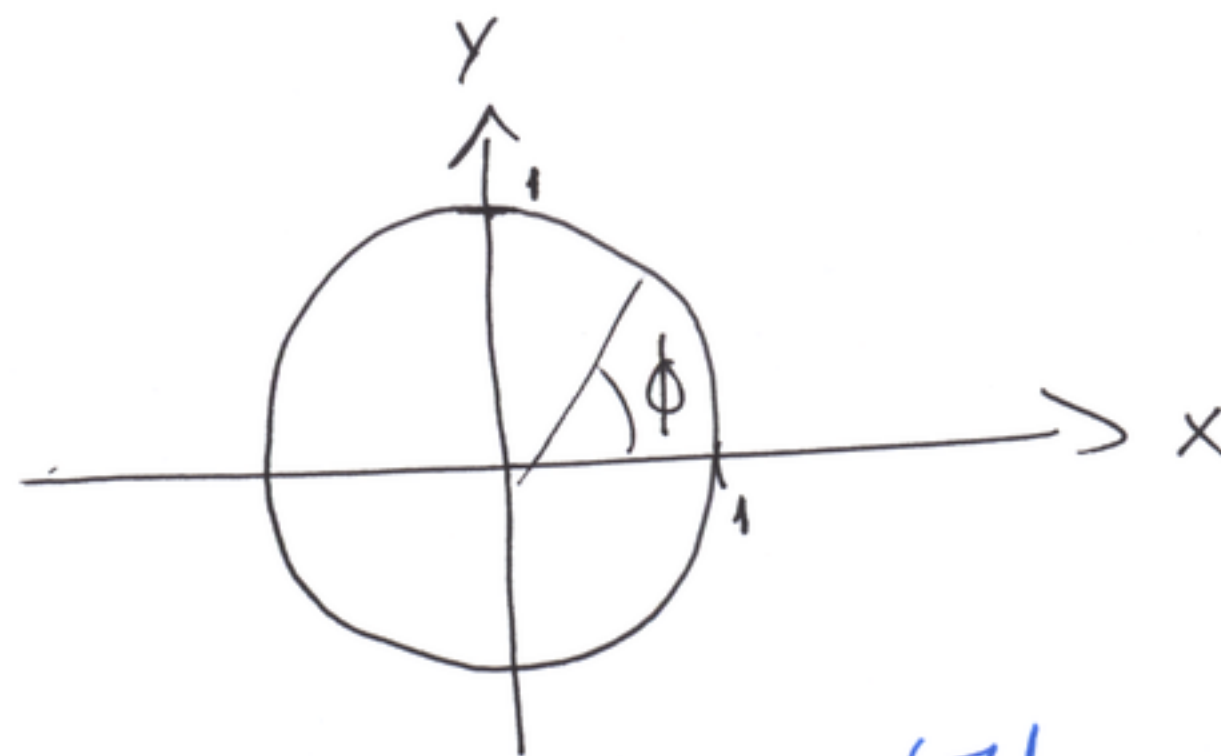
$$\text{normað } |E_{\theta}| = |\sin \theta|$$



lítil geislun í z-átt samsvöð
tvískautsvogi $\bar{P}$

(H-plane): Gefið $R, \theta = \pi/2$

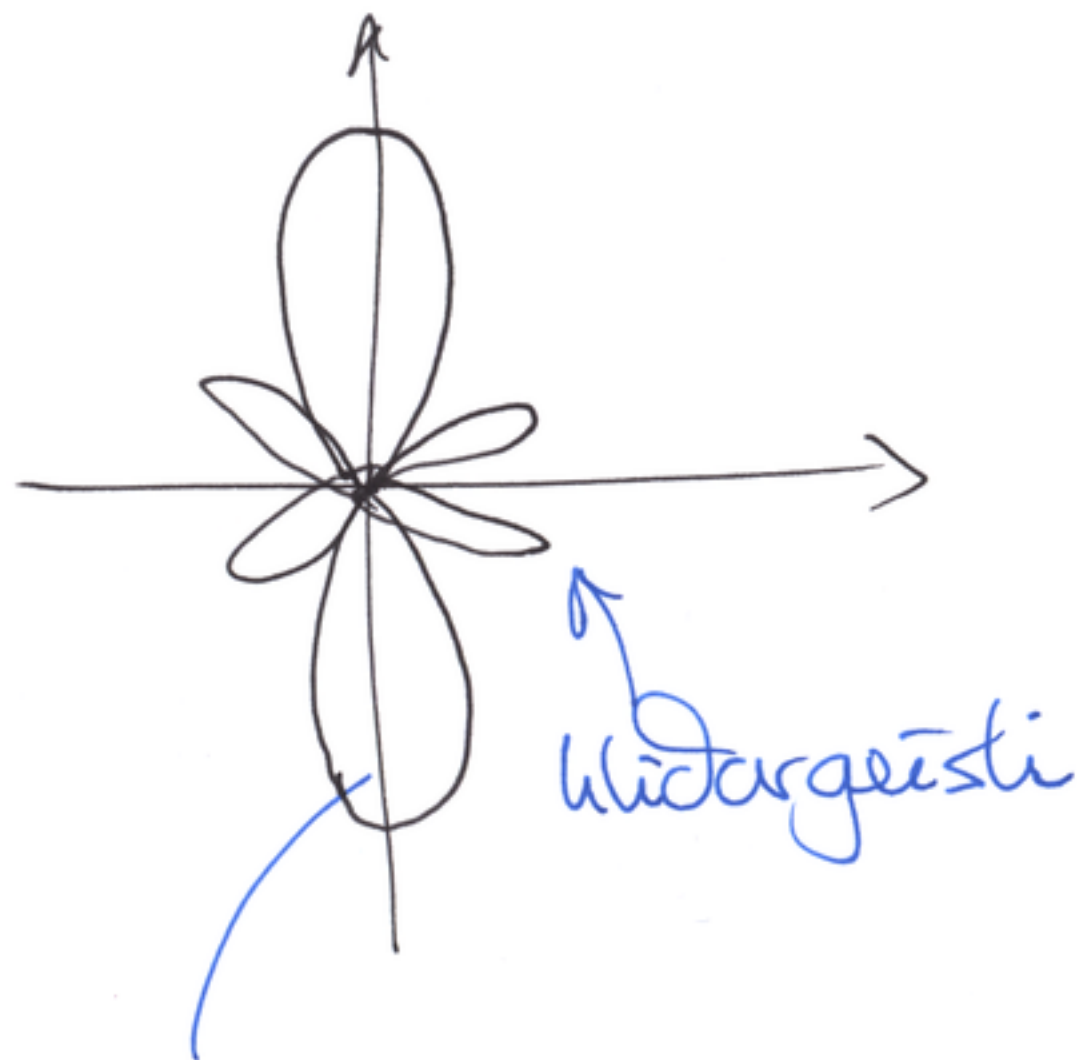
$$E_{\theta} = |\sin \theta| = 1$$



Jöfnu geislun í x-y-sléttu

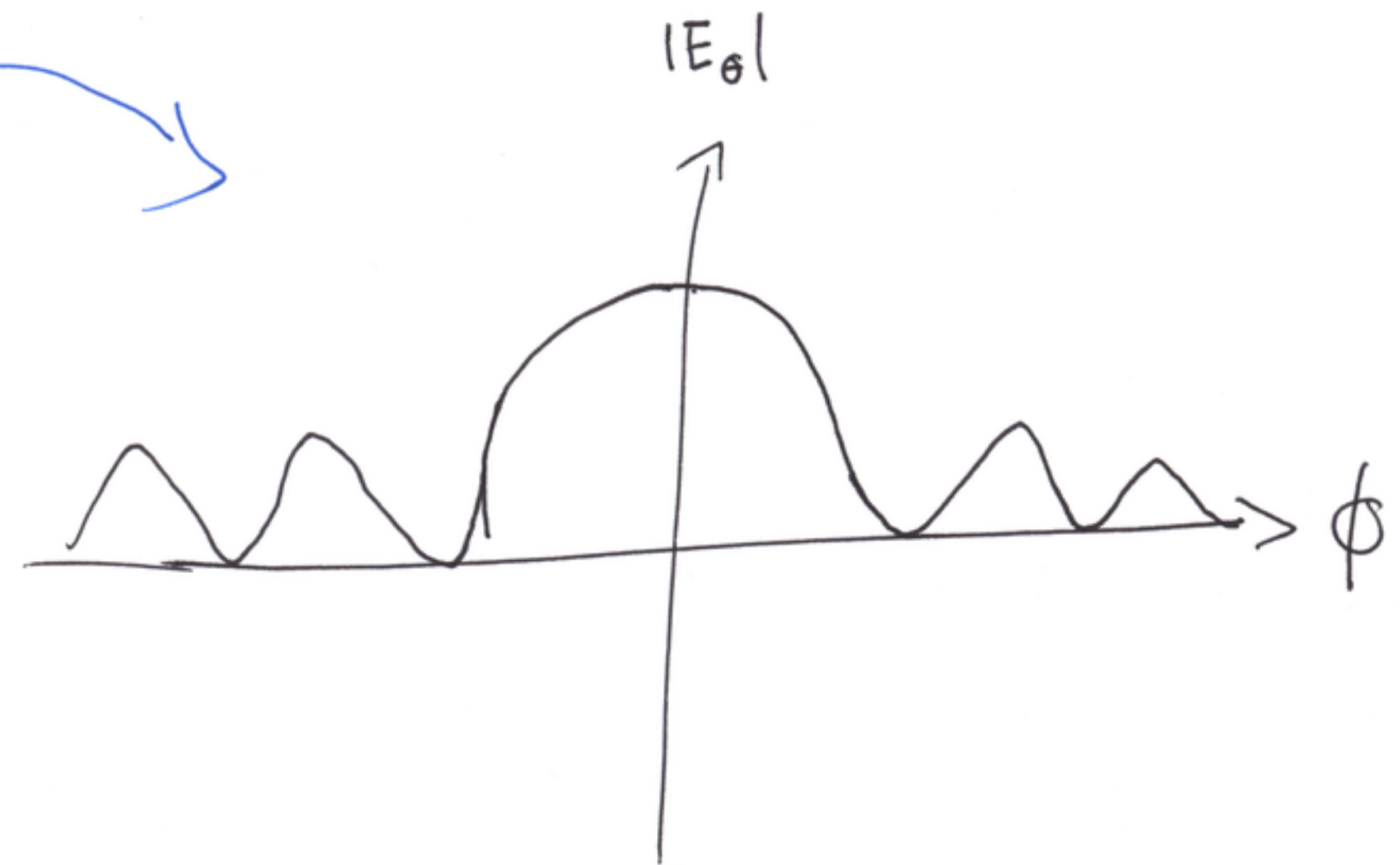
Sum loftnet geta höft
mynstur í x-y-slétta

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Adalgeisli

Önnur leið til
skotunar



Þá má skoða geisla breidd
og styrk hlutargeisla

Geislunarstyrkur, U

Er mældur í W/sr

rúmkorn, heil kúla er 4π (sr)

$$U = R^2 P_{ave}$$

og heildar afl geislað

$$P_r = \oint P_{ave} \cdot d\vec{s} = \oint U d\Omega$$

$$(d\Omega = \sin\theta d\theta d\phi)$$

Stefnumögnum er mæld með

$$G_D(\theta, \phi) = 4\pi \frac{U(\theta, \phi)}{P_r}$$

$$= \frac{4\pi U(\Omega)}{\oint U d\Omega}$$

Max Stefnumögnum er
áttun (directivity)

$$D = \frac{4\pi U_{max}}{P_r} \leftarrow$$

sem má skrifa sem

$$D = \frac{4\pi |E_{max}|^2}{\int |E(\Omega)|^2 d\Omega}$$

Geislað afl

Í loftnetinu og umhverfinu (jörðinni)
verður alltaf ómskt afltap P_e

Heildar inn-afl er því $P_i = P_r + P_e$

og aflmögnum loftnets er

$$G_P = \frac{4\pi U_{max}}{P_i}$$

Geislunarnýtni er skilgreind

$$\eta_r = \frac{G_P}{D} = \frac{P_r}{P_i}$$

Geislunar viðnám er gildi
þess ómska viðnáms sem eyðir
jafn miklu afli í hita og
loftnetið í geislu P_r



hátt geislunar viðnám
er hagkvæmt

Könnun þessa stíka með dæmi

Hertz-tvískaut

$$H_\phi = i \frac{Idl}{4\pi} \left(\frac{e^{-i\beta R}}{R} \right) \beta \sin\theta$$

$$E_\theta = i \frac{Idl}{4\pi} \left(\frac{e^{-i\beta R}}{R} \right) \eta_0 \beta \sin\theta$$

$$S_{ave} = \frac{1}{2} \Re |\vec{E} \times \vec{H}^*| = \frac{1}{2} |E_\theta| |H_\phi|$$

$$U = R^2 S_{ave} = \frac{(Idl)^2}{32\pi^2} \eta_0 \beta^2 \sin^2\theta$$

$$P_r = \oint U(\Omega) d\Omega = 2\pi \int_0^\pi U(\theta) \sin\theta d\theta$$

(5)

$$P_r = 2\pi \frac{(Idl)^2}{32\pi^2} \eta_0 \beta^2 \underbrace{\int_0^\pi \sin^3\theta d\theta}_{\left(\frac{1}{3} \cos^3\theta - \cos\theta \right) \Big|_0^\pi}$$
$$= -\frac{1}{3} + 1$$
$$-\frac{1}{3} + 1 = \frac{2}{3}$$

$$P_r = \frac{(Idl)^2}{12\pi} \eta_0 \beta^2$$

$$G_D(\Omega) = \frac{U(\Omega) 4\pi}{P_r}$$
$$= \frac{(Idl)^2 \eta_0 \beta^2 \sin^2\theta 12\pi}{8\pi (Idl)^2 \eta_0 \beta^2}$$
$$= \frac{3}{2} \sin^2\theta$$

$$G_D(\Omega) = \frac{3}{2} \sin^2 \theta$$

Mest geislu í x-y-slétta
og engin í pól áttirnar
tvar $\pm \hat{z}$

$$D = G_D\left(\frac{\pi}{2}, \phi\right) = 1.5$$

$$P_r = \frac{(I dl)^2}{12\pi} \eta_0 \beta^2$$

notum (8-14) með

$$\eta_0 \approx 120\pi$$

$$\beta = \frac{2\pi}{\lambda}$$

$$P_r = \frac{I^2}{2} \left\{ 80\pi^2 \left(\frac{dl}{\lambda} \right)^2 \right\}$$

til þess að finna geislunar-
viðnám notum við

$$P_r = \frac{1}{2} I^2 R_r$$

$$\rightarrow R_r = 80\pi \left(\frac{dl}{\lambda} \right)^2$$

Hér þarf að muna að sviðin okkar eru
aðeins rétt reiknað fyrir $dl \ll \lambda$



fækar lægt gildi fyrir R_r

Skodum virbút með löðni ∇

geisla a , lengd d sem Hertz-tístant

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Óvskt tap $P_e = \frac{1}{2} I^2 R_e$

R_e verður að tengja við yfirborsviðnám R_s

$$R_e = R_s \left(\frac{d}{2\pi a} \right) \quad \text{p.s.} \quad R_s = \sqrt{\frac{\pi f \mu_0}{\nabla}}$$

← frá flutningstímanum
sem við höfum
slépt

$$\eta_r = \frac{P_r}{P_r + P_e} = \frac{R_r}{R_r + R_e} = \frac{1}{1 + \left(\frac{R_e}{R_r} \right)}$$

$$\rightarrow \eta_r = \frac{1}{1 + \frac{R_s}{160\pi^3} \left(\frac{\lambda}{a} \right) \left(\frac{\lambda}{d} \right)}$$

Ef $\bar{u}$ $a = 1,8 \text{ mm}$, $d = 2 \text{ m}$, $f = 1,5 \text{ MHz}$

$\nabla = 5,8 \cdot 10^7 \text{ S/m}$ fyrir kopar fást $\lambda = 200 \text{ m}$

og $\eta_r = 0,58 \rightarrow \underline{58\% \text{ útyúi}}$

Jafnan

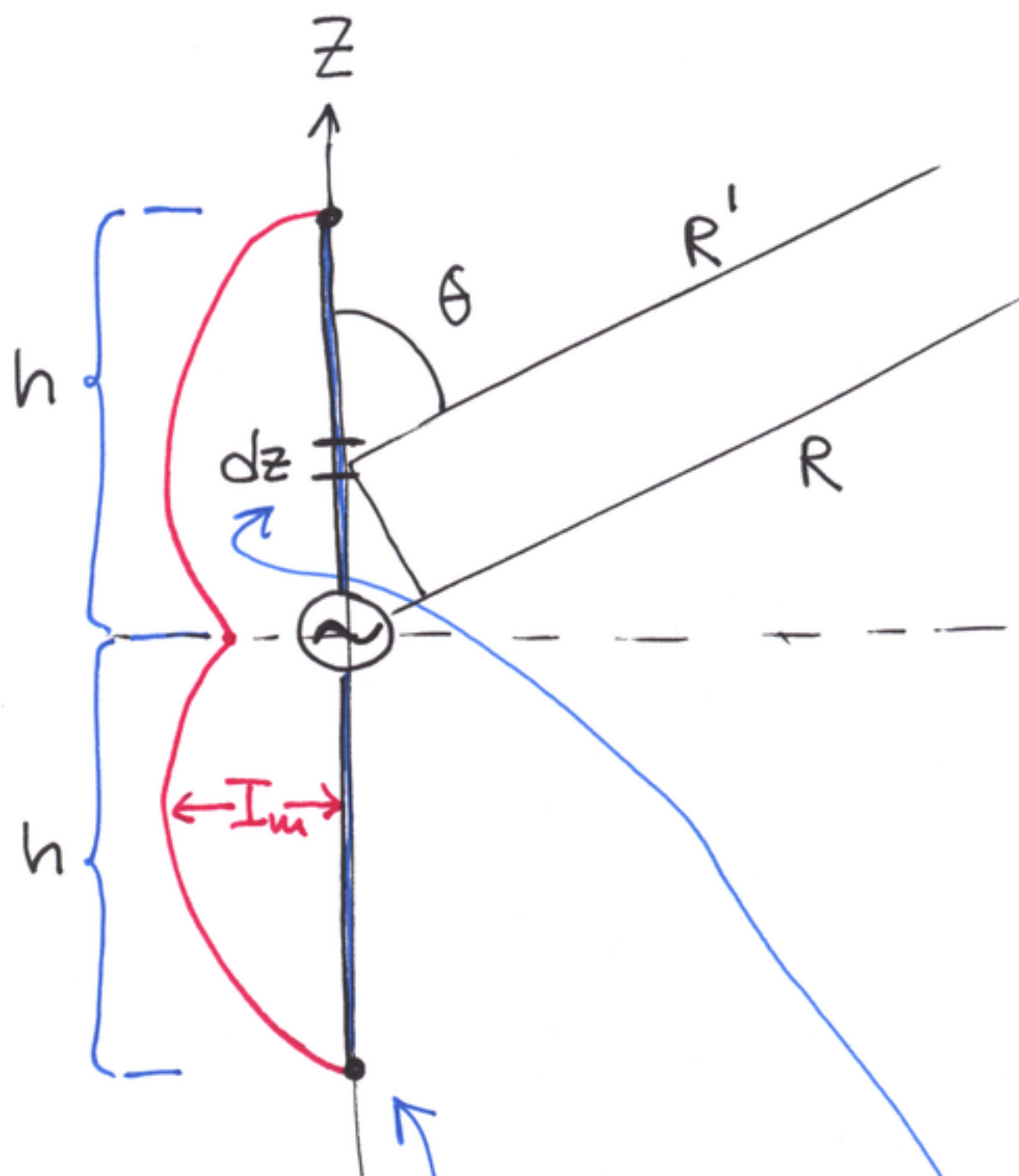
$$\eta_r = \frac{1}{1 + \frac{R_s}{160\pi^3} \left(\frac{\lambda}{a}\right) \left(\frac{\lambda}{d}\right)}$$

Sýnir að lág gildi fyrir $\left(\frac{a}{\lambda}\right)$ og $\left(\frac{d}{\lambda}\right)$

lækka útyúina (En jöfnur gilda aðeins í þessu merkgildi)

því þarftum við að stöðva loftnet með lengd sambærilega við bylgjulengdina λ

Grönn lína



Stráumurinn verður að hverfa í endanum

slæppum þú æð ákvörðaða stráuminn sjálf-samkvæmt og gerum ræð fyrir

$$I(z) = I_m \sin\{\beta(h - |z|)\}$$

$$= \begin{cases} I_m \sin\{\beta(h - z)\} & z > 0 \\ I_m \sin\{\beta(h + z)\} & z < 0 \end{cases}$$

Við látum okkur hugja að kanna fjar-sviðin

$$dE_\theta = \eta_0 dH_\phi = i \frac{I dz}{4\pi} \left(\frac{e^{-i\beta R'}}{R'} \right) \eta_0 \beta \sin\theta$$

fyrir litla bít loftnetsins dz

Mjög fjarri loftnetinu $R \gg h$

(10)

$$R' = (R^2 + z^2 - 2Rz \cos \theta)^{1/2} \approx R - z \cos \theta, \quad \frac{1}{R} \sim \frac{1}{R'}$$

$$E_\theta = \eta_0 H_\phi = i \frac{I_m \eta_0 \beta \sin \theta}{4\pi R} e^{-i\beta R} \int_{-h}^h \sin\{\beta(h-|z|)\} e^{i\beta z \cos \theta} dz$$

$$= i \frac{60 I_m}{R} e^{-i\beta R} F(\theta)$$



$$F(\theta) = \frac{\cos\{\beta h \cos \theta\} - \cos(\beta h)}{\sin \theta}$$

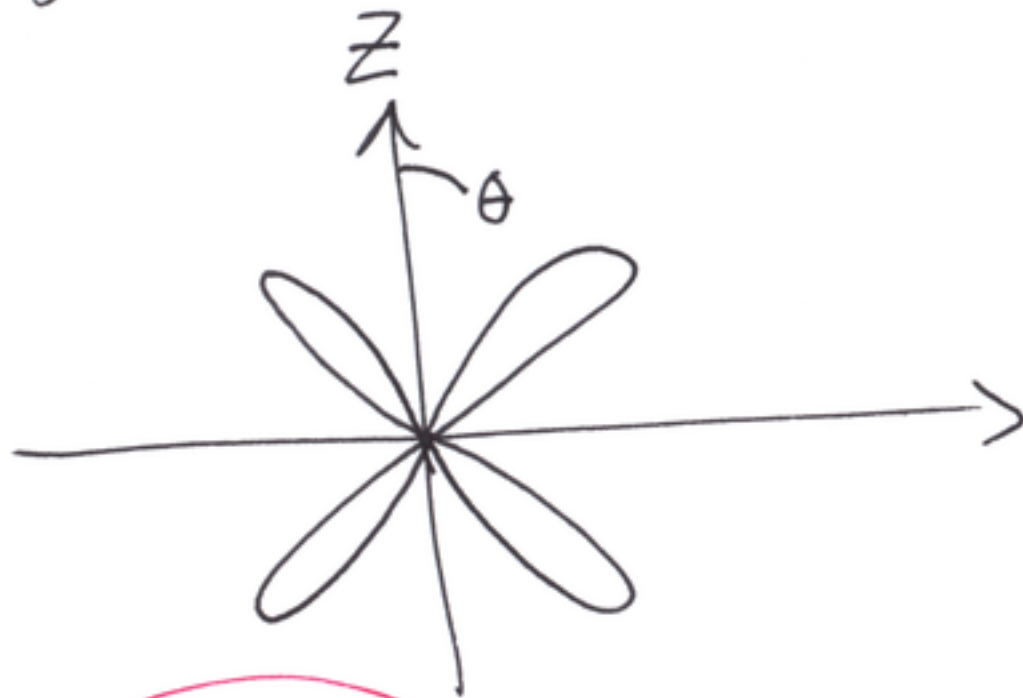


langd loftnetsins
skiptir öllumáli
hér

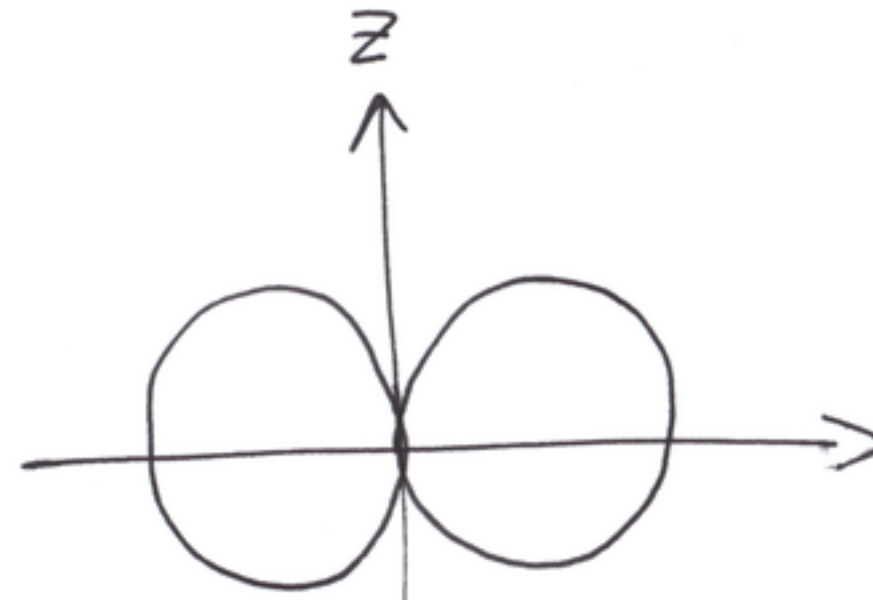
↗
E-plane myndast fallið fyrir þetta loftnet

fyrir $2h = 2\lambda$ er engin glæstur í $\pi/2$

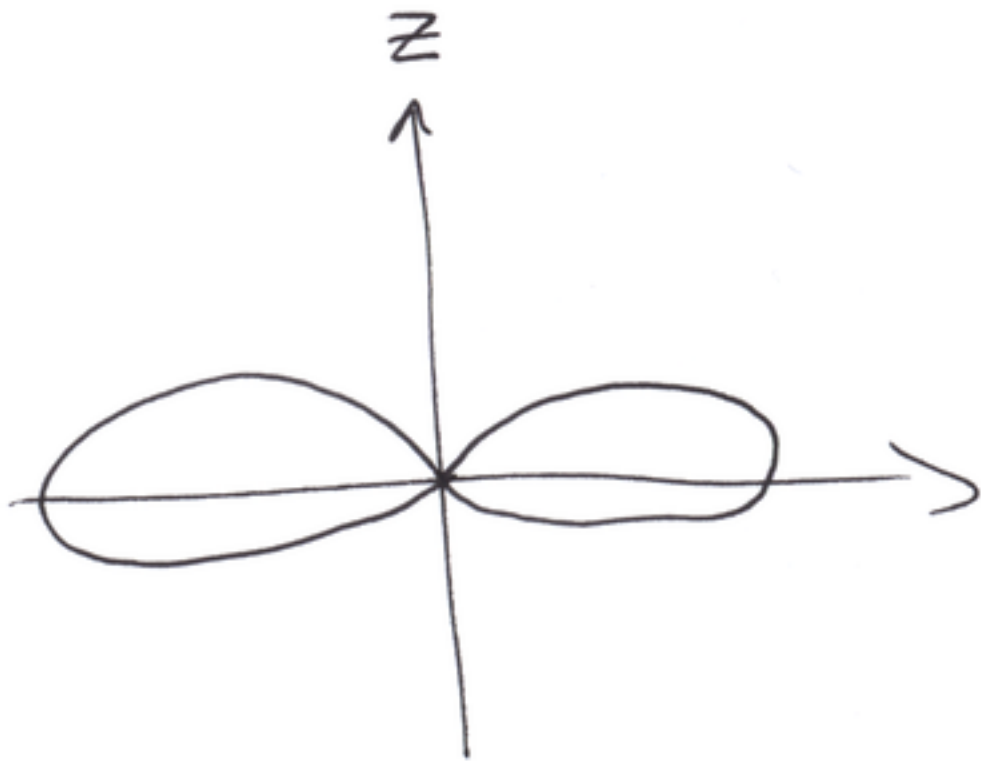
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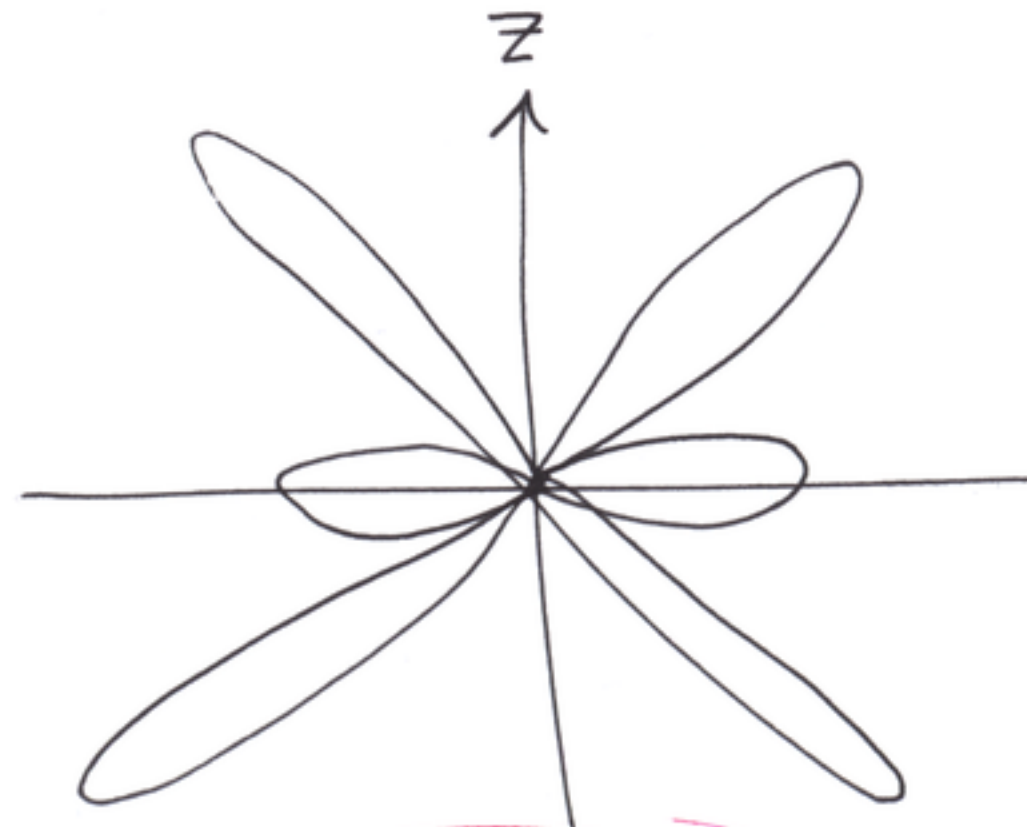
$$2h = 2\lambda$$



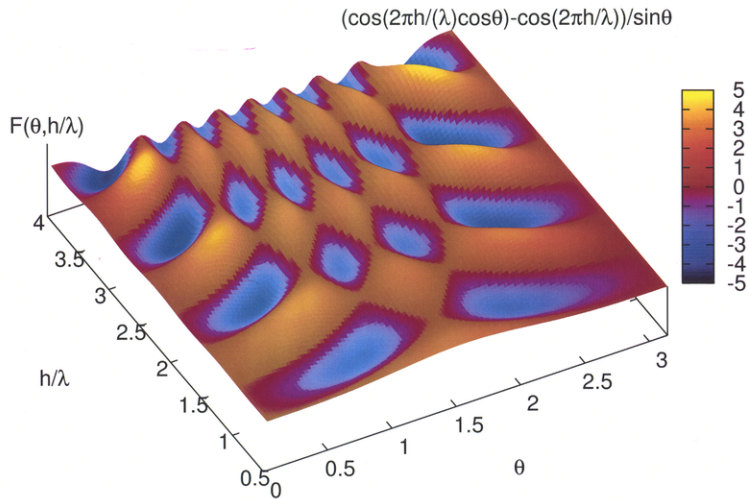
$$2h = \frac{1}{2}\lambda$$

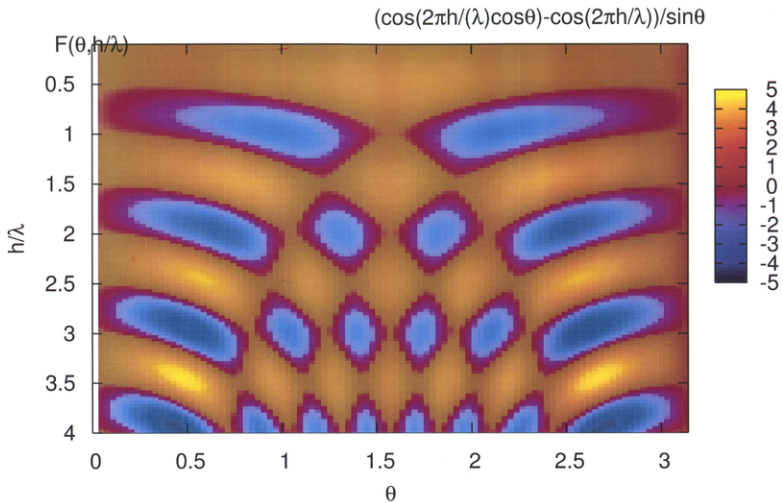


$$2h = \lambda$$

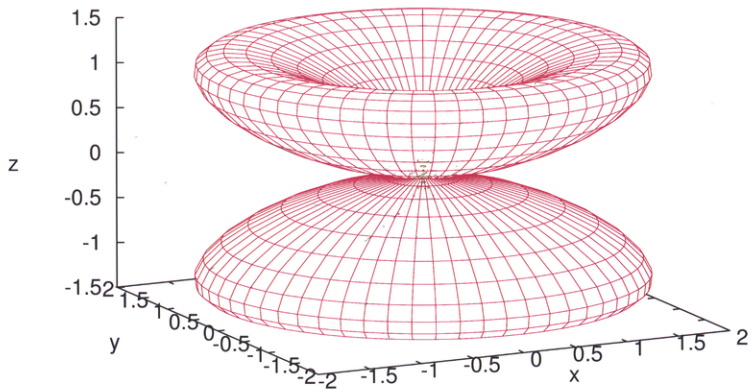


$$2h = \frac{3}{2}\lambda$$

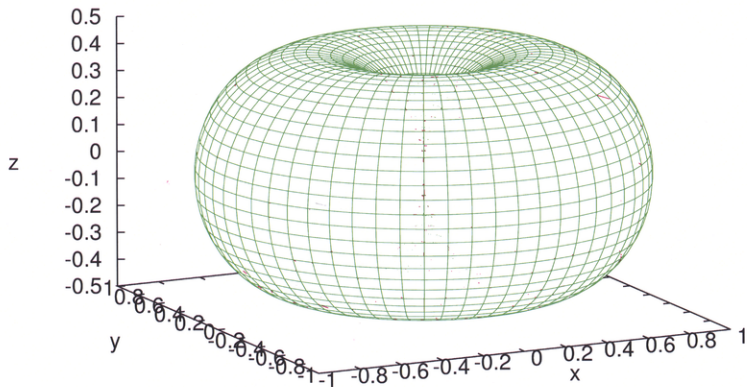




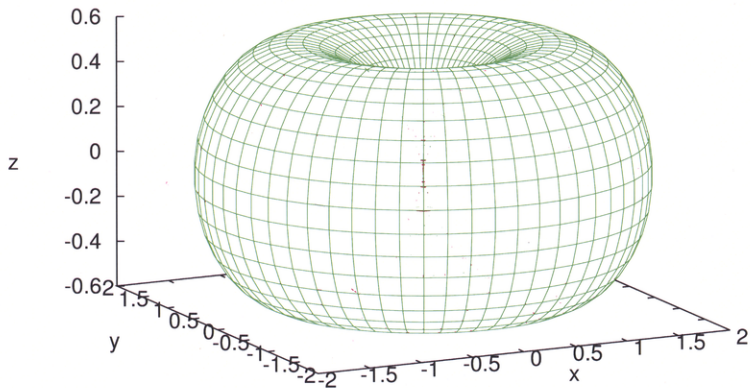
$$\beta\hbar=2\pi$$



$$\beta h = \pi/2$$



$$\beta h = \pi$$



$$\beta\hbar=3\pi/2$$

