

Problem 1

①

$$I = \int_V d\vec{r} e^{-\mu r} \left( \vec{\nabla} \cdot \frac{\hat{r}}{r^2} \right), \quad V: \text{sphere with radius } R$$

① use Eq. (1.99) which is the Gauss law in differential form for a point charge

$$\vec{\nabla} \cdot \left( \frac{\hat{r}}{r^2} \right) = 4\pi \delta^3(\vec{r})$$

$$\rightarrow I = \int_V d\vec{r} e^{-\mu r} 4\pi \delta^3(\vec{r}) = \underline{4\pi}$$

② use Eq. (1.59)

$$\int_V d\vec{r} f (\vec{\nabla} \cdot \vec{A}) = - \int_V d\vec{r} \vec{A} \cdot (\vec{\nabla} f) + \oint_S d\vec{a} \cdot f \vec{A}$$

$$\rightarrow I = - \int_V d\vec{r} \left( \frac{\hat{r}}{r^2} \right) \cdot (\vec{\nabla} e^{-\mu r}) + \oint_S d\vec{a} \cdot \left\{ \frac{e^{-\mu r} \hat{r}}{r^2} \right\}$$

$$= - \int_V d\vec{r} \left( \frac{\hat{r}}{r^2} \right) \cdot (-\mu e^{-\mu r} \hat{r}) + 4\pi e^{-\mu R}$$

$$\begin{aligned}
 &= \int_V d\tau \left( \mu \frac{e^{-\mu r}}{r^2} \right) + 4\pi e^{-\mu R} = \int_V r^2 dr d\Omega \mu \frac{e^{-\mu r}}{r^2} + 4\pi e^{-\mu R} \quad (2) \\
 &= \mu 4\pi \int_0^R dr (e^{-\mu r}) + 4\pi e^{-\mu R} = 4\pi \mu \frac{1}{\mu} \{1 - e^{-\mu R}\} + 4\pi e^{-\mu R} \\
 &= \underline{4\pi}
 \end{aligned}$$


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### Problem 3

$$(a) \quad \vec{F} = \underbrace{(\vec{r} \times \vec{p})}_{\text{vector}} \underbrace{(\vec{r} \cdot \vec{p})}_{\text{scalar}}$$

$$\text{use} \quad \vec{\nabla} \cdot (f \vec{A}) = f(\vec{\nabla} \cdot \vec{A}) + \vec{A} \cdot (\vec{\nabla} f)$$

$$\rightarrow \vec{\nabla} \cdot \vec{F} = (\vec{r} \cdot \vec{p}) \left[ \vec{\nabla} \cdot (\vec{r} \times \vec{p}) \right] + (\vec{r} \times \vec{p}) \cdot \left[ \vec{\nabla} (\vec{r} \cdot \vec{p}) \right]$$

and

$$\vec{\nabla} (\vec{r} \cdot \vec{p}) = \vec{p} \times (\vec{\nabla} \times \vec{r}) + (\vec{r} \cdot \vec{\nabla}) \vec{p} + (\vec{p} \cdot \vec{\nabla}) \vec{r} = (\vec{p} \cdot \vec{\nabla}) \vec{r}$$

$$\bar{\nabla} \cdot (\bar{r} \times \bar{\rho}) = \bar{\rho} \cdot (\bar{\nabla} \times \bar{r}) - \bar{r} \cdot (\bar{\nabla} \times \bar{\rho}) = 0$$

③

$$(\bar{\rho} \cdot \bar{\nabla}) \bar{r} = \bar{\rho}$$

$$\rightarrow \underline{\bar{\nabla} \cdot \bar{F} = (\bar{r} \times \bar{\rho}) \cdot \bar{\rho} = 0} \quad \text{as } \bar{r} \times \bar{\rho} \perp \bar{\rho}$$

$$\bar{\nabla} \times \bar{F}$$

$$\text{use } \bar{\nabla} \times (f \bar{A}) = f (\bar{\nabla} \times \bar{A}) - \bar{A} \times (\bar{\nabla} f)$$

$$\rightarrow \bar{\nabla} \times \bar{F} = (\bar{r} \cdot \bar{\rho}) \left\{ \bar{\nabla} \times (\bar{r} \times \bar{\rho}) \right\} - (\bar{r} \times \bar{\rho}) \times \left\{ \bar{\nabla} (\bar{r} \cdot \bar{\rho}) \right\}$$

$$\text{again } \bar{\nabla} (\bar{r} \cdot \bar{\rho}) = (\bar{\rho} \cdot \bar{\nabla}) \bar{r} = \bar{\rho}$$

$$\begin{aligned} \bar{\nabla} \times (\bar{r} \times \bar{\rho}) &= (\bar{\rho} \cdot \bar{\nabla}) \bar{r} - \underbrace{(\bar{r} \cdot \bar{\nabla}) \bar{\rho}}_0 + \underbrace{\bar{r} (\bar{\nabla} \cdot \bar{\rho})}_0 - \bar{\rho} \underbrace{(\bar{\nabla} \cdot \bar{r})}_3 \\ &= \bar{\rho} - 3\bar{\rho} = -2\bar{\rho} \end{aligned}$$

$$\rightarrow \bar{\nabla} \times \bar{F} = -(\bar{r} \cdot \bar{\rho}) 2\bar{\rho} - (\bar{r} \times \bar{\rho}) \times \bar{\rho}$$

$$= -2\vec{p}(\vec{r} \cdot \vec{p}) + \vec{p} \times (\vec{r} \times \vec{p})$$

$$= -2\vec{p}(\vec{r} \cdot \vec{p}) + p^2 \vec{r} - \vec{p}(\vec{p} \cdot \vec{r})$$

$$= \underline{-3\vec{p}(\vec{r} \cdot \vec{p}) + p^2 \vec{r}}$$

### Problem 2

$$\phi(\vec{r}) = q e^{-\mu r} / r, \quad r = |\vec{r}|$$

$\mu$ : constant

(a)  $\vec{E} = -\nabla \phi(\vec{r})$       3D spherical coordinates

$$\rightarrow \vec{E} = -q \frac{\partial}{\partial r} \left\{ \frac{e^{-\mu r}}{r} \right\} \hat{r}$$

$$= \hat{r} q e^{-\mu r} \left\{ \frac{\mu}{r} + \frac{1}{r^2} \right\}$$

Notice that  $\mu$  helps us to check the dimension

radial field  $\nearrow$

we notice that  $E$  dies off quicker than for a single point charge as  $r$  approaches infinity. This indicates that the total charge might be zero, but not all of it is in a single point.

(b) find  $\phi(\vec{r})$

From  $E$  we can conclude that it has a radial symmetry, spherical. we can "turn around" the differential form of the Gauss theorem

$$\begin{aligned}
 \phi(\vec{r}) &= \epsilon_0 \vec{\nabla} \cdot \vec{E} \\
 &= \epsilon_0 q \frac{1}{r^2} \frac{\partial}{\partial r} \left\{ e^{-\mu r} (\mu r + 1) \right\} \\
 &= \epsilon_0 q \frac{1}{r^2} \left\{ -\mu(\mu r + 1) + \mu \right\} e^{-\mu r} \\
 &= \epsilon_0 \frac{q}{r^2} \left\{ -\mu^2 r \right\} e^{-\mu r} = -\epsilon_0 \frac{q \mu^2}{r} e^{-\mu r}
 \end{aligned}$$

(6)

..., but check the "total charge" according to this distribution  $\rho(\vec{r})$

$$\begin{aligned}
 Q &= \int_V d\vec{r} \rho(\vec{r}) = -\epsilon_0 q \mu^2 \int_0^\infty r^2 dr d\Omega \frac{e^{-\mu r}}{r} \\
 &= -4\pi\epsilon_0 q \mu^2 \int_0^\infty r dr e^{-\mu r} = -4\pi\epsilon_0 q \mu^2 \frac{1}{\mu^2} \\
 &= -4\pi\epsilon_0 q
 \end{aligned}$$

which is negative  
if  $q > 0$

(c)

This does not allow the total E to die off exponentially as  $r \rightarrow \infty$

so we need to add a point charge at  $r = 0$ ,  $\rho$  could not be determined at 0 by our method, without a refinement.

$$\rho_{\text{Total}} = \rho(\vec{r}) + \epsilon_0 q \delta^3(\vec{r}) = \rho(\vec{r}) + \epsilon_0 q \frac{\delta(r)}{r^2}$$

This has relevance to a hydrogen atom in the ground state, for example....